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# THERMODYNAMICS OF THE Fm3m ≈ Pm3m TRANSITION IN THE POTASSIUM AND RUBIDIUM HALIDES\*

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Abstract – The Fm3m  $\rightleftharpoons$  Pm3m phase transition in the potassium and rubidium halides has been examined by a dilatometric method from 25°C to temperatures near the melting point. The pressure  $(P_{tr})$  and volume change  $(\Delta V_{tr})$  measured for these polymorphic transitions were treated by method of least squares. Averages of  $P_{tr}$  and  $\Delta V_{tr}$  at 25° were taken from this work and from data in the literature. These average values of  $P_{tr}$  and  $\Delta V_{tr}$  were used to calculate the  $P\Delta V$  work for this transition; the entropy change  $(\Delta S_{tr})$  was calculated from the average value of  $\Delta V_{tr}$  and the temperature coefficient,  $(dP_{tr}/dT)$ . Data are also given for  $P_{tr}$ ,  $\Delta V_{tr}$  and  $\Delta S_{tr}$  at the triple point (Fm3m  $\rightleftharpoons$  liquid  $\rightleftharpoons$  Pm3m). No solid-solid transition was observed in KF below 800°C and 45 kb. The following trends are noted. At constant pressure, the entropy change  $\Delta S$  (Fm3m  $\rightarrow$  Pm3m) is positive for the potassium salts KCl, KBr and K1; for the rubidium halides it is negative. The transition volume  $\Delta V$ (Fm3m  $\rightarrow$  Pm3m) in all cases decreases as the temperature increases.

#### 1. INTRODUCTION

THERE HAVE been numerous investigations of the pressure-induced polymorphic transition in the alkali halides since a transition of this type was first discovered in the rubidium halides RbBr and RbI by Slater[1] in 1926. Bridgman found a similar transition in RbCl [2] and in the potassium salts KCl, KBr and K1[3]. Piermarini and Weir[4,5] also reported the discovery of a solid-solid transition in the fluorides, RbF and KF. Jacobs[6] established, by high pressure X-ray diffraction, that the pressure-induced phase transition in RbI is from the sodium chloride type structure (space group Fm3m) to the denser cesium chloride type structure (space group Pm3m). Other workers [5, 7-9] have shown that this same type of structure change also occurs in RbCl, RbBr, KI and KCl. A recent report by Klement and Jayaraman[10] summarizes the work which has been carried out on the phase transitions in these salts.

Most of the measurements of the poly-

made at room temperature with the exception of the work by Bridgman[3], Pistorius[11, 12] and Pistorius and Snyman[13] at temperatures up to 200°C. We have studied the transition in the potassium and rubidium halides at higher temperatures by the volume discontinuity method[14, 15] to temperatures approaching the triple point between the two solid phases and the liquid phase. Data for the transition pressure ( $P_{tr}$ ), the transition volume ( $V_{tr}$ ) and their variation with temperature have been obtained and thermodynamic data were derived from these experimentally measured quantities. These will be presented and discussed in this paper.

morphic transition in these salts have been

#### 2. EXPERIMENTAL

### (A) Materials

The salts used in these measurements were Mallinckrodt analytical reagent grade with the exception of RbF which was obtained from the A. D. McKay Co. The fluorides. KF and RbF were dried at 300°C under vacuum ( $P < 10^{-6}$  Torr.) for 48 hr. This was followed by treatment with anhydrous HF at a pressure of 1 atm for 24 hr. The chlorides, bromides and iodides of these alkali metals were heated for 24 hr under high vacuum at a temperature of 150°C. The dried salts were compressed into pellets at a pressure of ~ 2 kb. The fluorides, KF and RbF, were

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handled in an atmosphere of argon in order to reduce pickup of water.

#### (B) Apparatus

Use of the piston displacement method for the detection of phase transitions has been described by Bridgman [14], Kennedy and Newton[15], and by Jayaraman *et al.* [16]. Our techniques are similar to those described by the above workers. However, we have extended the temperature range of this piston displacement method from the upper limit of  $\sim 200^{\circ}$ C, imposed by external heating of the compound cylinder apparatus, to temperatures of  $\sim 1000^{\circ}$ C which are attainable with use of the internally heated graphite furnace. In this modification the sample in the form of a cylinder 0.6 cm in dia. by 1.5 cm in length was placed within a cylindrical graphite furnace heating element within the high pressure chamber. The volume of the sample in this case now comprises approximately 20 per cent of the total volume of the pressure chamber.

The transition volume was determined by the piston displacement method. This technique has been used by Bridgman[14] and more recently described by Jayaraman *et al.* [17]. Usually the sample and piston have the same diameter; however, in the technique used here the sample and piston differ in diameter by a factor of two. Thus the sample and internal furnace apparatus must distort or 'flow' in order that the volume change which takes place in the sample be transmitted into a displacement of the piston. The greater reproducibility of the  $\Delta V_{tr}$  measurements at higher temperatures, i.e.  $t > 200^{\circ}$ C, is very probably due to the increase in the plasticity of the furnace parts and of the sample at these higher temperatures.

For pressure calibration the phase transition  $Bi(1) \rightleftharpoons Bi(11)$  and  $Bi(11) \rightleftharpoons Bi(11)$  at 25°C were used. We find 25.6±0.2 kb and 27.3±0.2 kb respectively for these transitions. Kennedy and LaMori[18, 19] give  $25.38 \pm 0.02$  kb and  $27.0 \pm 0.2$  kb for these transitions in Bi at 25°C. It was assumed that this calibration holds good at higher temperatures also.

Temperatures were measured with a chromel-alumel thermocouple whose junction was positioned in the geometric center of the salt sample. No corrections for the pressure effect on e.m.f. were made since corrections at pressures below 25 kb are expected to be small, i.e. less than 1% for Cr-Al thermocouples[20].

#### 3. RESULTS AND DISCUSSION

The pressures of this polymorphic transition ( $P_{tr}$ ) were determined upon compression and decompression at approximately 100° intervals, from 25 to 700°C and in some cases up to 800°C. The results for the salts KCl, KBr and KI are shown in Fig. 1; similarly, results for RbF, RbCl, RbBr and RbI are given in Fig. 2. No phase transition was ob-





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Fig. 2. *PV* diagram for the rubidium halides: □-melting point data of Clark (1959). ■-melting point data of Pistorius (1965, 1966), ○-Fm3m ≈ Pm3m transition, this work.

served in KF at pressures up to 45 kb and at temperatures up to 800°C.

As in the case of most solid-solid phase transitions[21] these salts also show a hysteresis effect. The width of the hysteresis loop is, in general, less than one kilobar at temperatures greater than 200°C, and becomes larger and less reproducible at temperatures below 200°C. We have assumed that the equilibrium transition pressure is the average of the transition pressures obtained on compression and decompression. In all cases, the equilibrium transition pressure appears, within the experimental accuracy of these data, to be a linear function of the temperature over the interval  $200-700^{\circ}$ C. These  $P_{tr}$  data were treated by method of least squares to obtain the transition pressure as a function of temperature in the form

$$P_{\rm tr}(\rm kb) = a + bt \tag{1}$$

where  $P_{tr}$  is the Fm3m  $\rightleftharpoons$  Pm3m transition pressure at a given temperature t (°C). Data for the constants a and b are given in Table 1. The constant b is of particular significance since b = dP/dT. The values for the transition pressure at 25°C reported in Table 2 were obtained from extrapolation of the high temperature data, by the use of equation (1).

The transition volumes\*  $(\Delta V_{tr}, Fm3m \rightarrow Pm3m)$  for KCl, KBr and KI at temperatures between 200-800°C are shown in Fig. 3. Similarly,  $\Delta V_{tr}$  for RbF, RbCl, RbBr and RbI are shown in Fig. 4. The volume change  $(\Delta V_{tr})$  appears to be a linear function of temperature at temperatures above 200°C. These data were also treated by method of least squares to obtain the transition volume as a function of temperature in the form

$$\Delta V_{\rm tr}(\rm cm^3/mole) = c + dt.$$
 (2)

where  $\Delta V_{tr}$  is the transition volume at a given temperature t (°C). Values for the constants c and d are given in Table 1. The transition volumes at 25°C reported in Table 3 were obtained from an extrapolation of the high tem-

\*The transition volume vs. temperature data shown in Figs. 3 and 4 are not at constant pressure (except in the special case of KCl where  $dT/dP \approx 0$ ) but are at pressures corresponding to the temperatures along the Fm3m  $\rightleftharpoons$  Pm3m phase boundary.



Fig. 3.  $\Delta V(Fm3m \rightarrow Pm3m)$  vs. T for KCl, KBr and Kl.

Table 1. Constants for the temperature dependence of the Fm3m  $\rightleftharpoons$  Pm3m transition pressure  $[P_{tr}(kb) = a + bt]$  and the transition volume Fm3m  $\rightarrow$  Pm3m  $[\Delta V_{tr}(cm^3/mole) = c + dt]$  in the potassium and rubidium halides

$d \times 10^3$
(cms/mole/deg)
$0.830 \pm 0.176$
$0.869 \pm 0.201$
$1 \cdot 211 \pm 0 \cdot 295$
$0.445 \pm 0.513$
$2.880 \pm 0.269$
$1.597 \pm 0.366$
$1.920\pm0.236$

Salt	Bridgman 1945[25]	Daniels. <i>et al.</i> 1966[31]	Kennedy and LaMori, 1962[18]	Knof and Maisch. 1963[26]	Larson, 1965[29]	Piermarini and Weir, 1962[4]	Pistorius, 1964[11, 12]	Pistorius and Snyman. 1964[13]	This work	Average Value
									No	ne
KF						35		$14.6 \pm 1.3$	obsei	rved
KC1	19.7		$18.28 \pm 0.21$		20.0		$19.27 \pm 0.08$		$19.55 \pm 0.13$	$19.4 \pm 0.3$
KBr	18.0		$17.88 \pm 0.06$		18.5		$17.43 \pm 0.07$		$17.99 \pm 0.28$	$18.0 \pm 0.2$
KI	17.8		$17.48 \pm 0.24$				$17.34 \pm 0.05$		$18.27 \pm 0.19$	$17.7 \pm 0.3$
RbF				33		12		6.1	$34.47 \pm 0.38$	$33.8 \pm 0.8$
RbC1	4.90						5.28		$5.68 \pm 0.14$	$5.32 \pm 0.24$
RbBr	4.50						$4.20 \pm 0.2$		$4.92 \pm 0.06$	$4.57 \pm 0.35$
Rbl	3.96	3.54					$3.59 \pm 0.16$		$3.68 \pm 0.05$	$3.69 \pm 0.11$

Table 2. Transition pressure ( $P_{tr}$ , kb) for the Fm3m  $\rightleftharpoons$  Pm3m phase change in the potassium and rubidium halides at 25°C

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Table	3.	Transition	volume	$(-\Delta V_{\rm tr},$	cm <sup>3</sup> /mole)	for the	$Fm3m \rightarrow$	Pm3m	phase	change	in th	e potassiu	n and
					rubia	lium hali	ides at 25°C	2					

Salt	Adams and Davis, 1962[8]	Bridgman. 1945[25]	Genshaft <i>et al.</i> , 1967[30]	Jacobs, 1938[6]	Jamison, 1957[7]	Nagasaki and Minomura, 1964[9]	Pistorius and Snyman, 1964[13]	Weir and Piermarini, 1964[5]	This work	Average value
KF							1.0	2.49	Non	e Obs.
KC1		4.20				3.85		6.85	$4 \cdot 11 \pm 0 \cdot 10$	$4.05 \pm 0.17$
KBr		4.55						8.35	$4 \cdot 17 \pm 0 \cdot 11$	$4.36 \pm 0.19$
KI		4.50			4.50			11.8	$4{\cdot}41\pm0{\cdot}15$	$4{\cdot}47\pm0{\cdot}13$
RbF								3.70	$1.83 \pm 0.29$	
RbCl	6.55	6.00	6.30					5.76	$6.95 \pm 0.11$	$6.30 \pm 0.35$
RbBr		6.55						6.60	$7.43 \pm 0.18$	$6.86 \pm 0.39$
RbI	9.65	7.50		7.9				8.26	$8{\cdot}10\pm0{\cdot}10$	$8{\cdot}28\pm0{\cdot}31$

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Fig. 4.  $\Delta V(Fm3m \rightarrow Pm3m)$  vs. T for the rubidium halides.

perature data by use of equation (2). The pressure-temperature coordinates of the Fm3m  $\rightleftharpoons$  liquid  $\rightleftharpoons$  Pm3m triple point for these salts were determined from the intersection of the melting curves of the Fm3m and Pm3m phases at the Fm3m  $\rightleftharpoons$  Pm3m phase boundary. The melting curves of the Fm3m and Pm3m phases were drawn through the data points of Clark[22] and of Pistorius [23] in order that the intersection of these two curves would fall on the Fm3m  $\rightleftharpoons$  Pm3m phase boundry obtained in this work. Triple point coordinates obtained for the salts KCl, KBr and K1 are shown in Fig. 1; coordinates for the rubidium halides are shown in Fig. 2. These triple point P-T coordinates are compared with the P-T coordinates given by Clark[22] and by Pistorius[23] in Table 4.

## Potassium halides

Potassium fluoride was examined at pressures up to 45 kb at approx.  $100^{\circ}$  intervals from room temperature up to  $800^{\circ}$ C. However, we fail to find the phase transition reported by Weir and Piermarini[5] and by Pistorius and Snyman[13]. Pistorius *et al.* found the volume change of this transition to be small, i.e. 0.5 per cent. The sensitivity of the method used here is more than adequate to detect a phase transition with such a small

Salt	Clark[22]		Pistorius[23]		This work							
	P (kb)	Т (°С)	P (kb)	<i>T</i> (°C)	P (kb)	<i>Т</i> (°С)	$-\Delta V_{\rm tr}$ (cm <sup>3</sup> /mole)	$\Delta S_{\rm tr}$ (cal/mole/deg)	$T \Delta S$ (cal/mole)	P (cal/mole)	$\Delta E$ (cal/mole	
KF					none o	bserved						
KC1	$18.95 \pm 0.1$	$1042 \pm 5$	$18.7 \pm 0.7$	$1050\pm10$	$19.5 \pm 0.3$	$1070 \pm 10$	$3.24 \pm 0.21$	+0.001	17	-1520	-1540	
KBr			$16.8 \pm 1.0$	$997 \pm 10$	$16.8 \pm 0.6$	$1005 \pm 10$	$3.32 \pm 0.18$	+ 0.09	120	-1340	-1460	
KI			$16{\cdot}9\pm1{\cdot}5$	$934\pm15$	$15{\cdot}3\pm0{\cdot}4$	$940\pm15$	$3 \cdot 31 \pm 0 \cdot 19$	+0.25	308	-1210	-1520	
RbF			32	1055	$38.5\pm0.8$	$1105 \pm 15$	$1.39 \pm 0.55$	-0.12	-170	-1280	-1110	
<b>RbCl</b>	7.80	852	$7.80 \pm 0.01$	$852 \pm 5$	$8 \cdot 1 \pm 0 \cdot 3$	$870 \pm 5$	$4.45 \pm 0.18$	-0.30	-344	-860	-520	
RbBr			$6.1 \pm 0.8$	$808 \pm 10$	$6 \cdot 0 \pm 0 \cdot 2$	$830 \pm 10$	$6.15 \pm 0.23$	-0.20	-222	-880	-660	
RbI			$5.0 \pm 0.8$	$760 \pm 10$	$4.7\pm0.2$	$763 \pm 10$	$6.69 \pm 0.16$	-0.22	-233	-750	-520	

Table 4. Thermodynamic properties of the $Fm3m \rightarrow Pm3m$ transition in the potassium and rubidium halides at the
$Fm3m \rightleftharpoons liquid \rightleftharpoons Pm3m triple point$

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volume change. Since we find no polymorphic change at high temperatures where the transition kinetics would be more favorable, and the fact that Pistorius[23] found no break in the melting curve suggests that KF has no solidsolid transition at pressures below 45 kb.

The Fm3m  $\rightleftharpoons$  Pm3m transition pressures found for the salts KCl, KBr and Kl are compared in Table 2 with the transition pressures reported in the literature. Bridgman has examined this transition in these salts on three separate occasions [3, 24, 25]; however, only the latest of Bridgman's data [25] are given. In general our results for  $P_{tr}$  are in good agreement with the published data (Table 2). Average values of  $P_{tr}$  were in each case calculated from the data tabulated in Table 2.

Bridgman[3] and Pistorius[11, 12] have examined the effect of temperature on the transition pressure in these salts over the temperature range 25°-200°C. Bridgman[3] gives +3.45, -4.54 and +3.33 bar/deg respectively for the temperature dependence of the transition pressure in KCl. KBr and KI. Pistorius [11, 12], on the other hand, gives -0.25, +0.55and -1.88 bar/deg for these same salts. These data are compared with the constant b from Table 1, i.e. the value of dP/dT found in this work. The data from these three sources all indicate that the variation of  $P_{\rm tr}$  with temperature is small. It is seen that the values of dP/dT found by Bridgman and Pistorius all disagree in sign. The sign of dP/dT however determines the sign of the entropy change,  $\Delta S_{\rm tr}$  (Fm3m  $\rightarrow$  Pm3m) for this transition in these salts since  $\Delta V_{\rm tr}$  (Fm3m  $\rightarrow$  Pm3m) is negative in all cases.

The agreement between  $\Delta V_{tr}$  for KCl, KBr and KI found here and with published data (Table 3) is also relatively good, with the exception of the  $\Delta V_{tr}$  data given by Weir and Piermarini[5]. Weir *et al.* obtained their transition volume data from high pressure X-ray diffraction studies. Their data, however, are also in disagreement with the  $\Delta V_{tr}$  data given by Jamieson[7] and by Nagasaki *et al.*[9] for KI and KCl respectively. (The latter data for KI and KCl were also determined from X-ray diffraction studies carried out at high pressures.) These  $\Delta V_{tr}$  data of Weir *et al.*[5] were not used in the calculation of the average value of  $\Delta V_{tr}$  for the salts KCl, KBr and KI.

## Rubidium halides

Piermarini and Weir[4] report a solidsolid phase transaction in RbF at a pressure between 9 and 15 kb. They give the volume change for this transition as -20 per cent. Knof and Maisch[26] have observed changes in the optical transmission properties of RbF at a pressure of 33 kb, which is similar to the optical effect associated with the Fm3m  $\rightarrow$ Pm3m transition in the potassium halides, KCl and KBr. Pistorius and Snyman[13], on the other hand, report phase transitions in RbF at a pressure of 6.1 kb at a temperature of 20°, and 5.0 kb at 200°C. The latter authors[13] noted that this phase transition in RbF was very sluggish. No indication of the size of  $\Delta V$  for this transition was given by Pistorius et al.[13].

The results from our dilatometric measurements clearly indicate a polymorphic transition in RbF at about 35 kb, in agreement with the results reported by Knof et al. [26]. Transition pressure data for RbF are shown in Fig. 2. The transition was very sluggish at 100° and no transition was observed at room temperatures. However, at temperatures above 300°C the transition proceeds as readily as was found in the other alkali halides at 200°C. In the case of RbF, data taken only above 300°C were used in the least squares treatment of  $P_{\rm tr}$  and  $\Delta V_{\rm tr}$  vs. temperature. It was also necessary to modify the experimental arrangement since at temperatures above 400°C RbF reacts with the graphite heater causing it to become relatively non-conducting. Pistorius[23] found Ni to be a satisfactory container for RbF at high temperatures. We therefore isolated our sample from the graphite heater with a thin (0.0025 cm) nickel sleeve. With this modification our measurements were still limited to temperatures below about

750°C since an irreversible change occurs in the resistance of graphite at higher temperatures, presumably from attack by RbF through the nickel foil. Our RbF samples were examined by X-ray diffraction after being removed from the pressure chamber and showed diffraction lines characteristic of only the Fm3m phase. Data for  $P_{\rm tr}$  and  $\Delta V_{\rm tr}$  for this salt at 25°C are given in Tables 2 and 3 respectively.

The pressures obtained for the Fm3m  $\rightleftharpoons$  Pm3m transition in the rubidium halides are compared with the published values in Table 2. Averages of  $P_{tr}$  for each of these salts were calculated from data given in Table 2. However, in the case of RbF only the data of Knof and Maisch[26] were averaged with the data from this work. Pistorius[11, 12] gives the temperature dependence of the polymorphic transition pressure in RbCl, RbBr and RbI as 3.01, 2.19 and 1.94 bar/deg respectively. The agreement with dP/dT found here (*b* in Table 1) is better than in the case of the potassium salts.

### Thermodynamic calculations

The average values of  $P_{tr}$  and  $\Delta V_{tr}$  given in Tables 2 and 3 were used to calculate the thermodynamic properties of this polymorphic transition in these salts at a temperature of 298°K. However in the case of dP/dT, we have used only the values obtained in this work, since our measurements are over a much wider temperature interval (from 400°C to as much as 700°C) than were Bridgman's [24] and Pistorius'[12, 13] work.

At a temperature and pressure where the two polymorphs may be considered to be in equilibrium, then

$$G(Fm3m phase) = G(Pm3m phase).$$
 (3)

Thus

$$\Delta G = 0 \tag{4}$$

and at constant temperature and pressure

$$\Delta E = P \Delta V - T \Delta S. \tag{5}$$

The external  $P\Delta V$  work necessary to carry out the Fm3m  $\rightarrow$  Pm3m phase transformation (reversibly) at constant pressure and temperature is calculated from the average values of  $P_{tr}$  and  $\Delta V_{tr}$ . These data for  $P\Delta V$  at 298°K are given in Table 5. Data for  $P\Delta V$  reported by Weir *et al.*[5] are also shown. The larger values of  $P\Delta V$  reported by these authors[5] for the potassium salts arise from the larger transition volumes found by them. On the other hand, good agreement in case of RbF is fortuitous since Weir's much lower value of  $P_{tr}$  is compensated for by the larger value of

Table 5. Thermodynamic properties of the  $Fm3m \rightarrow Pm3m$  transition in the potassium and rubidium halides at 298°K

Salt	$\Delta S_{\rm tr}$ Pistorius[11, 12] (cal/mole/deg)	$\Delta S_{tr}^*$ This work (cal/mole/deg)	$T\Delta S$ This work (cal/mole)	$P\Delta V$ Weir <i>et al.</i> [5] (cal/mole)	$P\Delta V^{\dagger}$ This work (cal/mole)	$\Delta E$ This work (cal/mole)
KCI	+0.03	0.002	6	-3.000	-1880	-1890
KBr	-0.11	0.12	37	-3,300	- 1880	-1920
KI	+0.53	0.34	103	-4,500	-1890	- 1990
RbF		-0.16	-48	-1,300	1490	-1440
RbC1	-0.44	-0.42	-125	-700	-800	-675
RbBr	-0.35	-0.23	-67	-700	-750	-680
Rbl	-0.38	-0.28	- 83	-700	-730	-660

\* Average values of  $\Delta V_{tr}$  in Table 4 were used to calculate these entropies.

†Average value of  $P_{tr}$  in Table 1 and average value of  $\Delta V_{tr}$  in Table 4 were used to calculate this  $P\Delta V$  product.

 $\Delta V$ . The agreement of  $P\Delta V$  from this work and from Weir's is, however, relatively good in the case of RbCl, RbBr and RbI.

The entropy change ( $\Delta S_{tr}$ , Fm3m  $\rightarrow$  Pm3m) was calculated from the Clausius Clapeyron relation. Data obtained for  $\Delta S_{tr}$  at 298°K are given in Table 5. Similarly, Pistorius[12, 13] has calculated  $\Delta H$  for this polymorphic transition in these salts from his value of dP/dT and from Weir's and Piermarini's [5] transition volume. As noted above, Weir's value of  $\Delta V_{\rm tr}$  for the potassium salts are considerably higher than the results found by others (results, on the other hand, for the rubidium salts RbCl, RbBr and RbI are in good agreement, Table 2). The entropy of transformation derived from Pistorius' calculations are -0.44, -0.35 and -0.38 cal/mole/ deg respectively for RbCl, RbBr and RbI. The agreement between the  $\Delta S_{tr}$  calculated by Pistorius[12, 13] and those derived from this work is relatively good for these salts.

Data for the  $T\Delta S$  term in equation (5) are given in Table 5 as also the data for the change in internal energy ( $\Delta E$ ) associated with this polymorphic transformation. Similarly, thermodynamic properties for this Fm3m  $\rightarrow$ Pm3m transformation at the triple point are given in Table 4.

Several trends in the transition properties of these salts are noted. From Figs. 4 and 5, it is seen that in all cases studied, the absolute value of transition volume decreases with an increase in temperature. In general, the pressure of the Fm3m  $\rightleftharpoons$  Pm3m transition is relatively insensitive to temperature for the halides of these two alkali metals. The potassium salts do, however, show a small decrease whereas the rubidium salts show a correspondingly small increase in transition pressure with an increase in temperature. The transition pressure for CsCl is also known to increase with an increase in temperature [27]. Consequently, at constant pressure  $\Delta S$ - $(Fm3m \rightarrow Pm3m)$  is positive for the potassium halides and negative for the rubidium halides (and CsCl). Thus in the case of the alkali chlorides, the absolute entropy of the Fm3m phase becomes larger relative to the Pm3m phase as the cation increases in atomic number and therefore in size. A similar trend has been noted [28] in the fusion of the Fm3m phase of these chlorides.

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#### REFERENCES

- 1. SLATER J. C., Proc. Am. Acad. Arts Sci. 61, 144 (1926).
- 2. BRIDGMAN P. W., Z. Krist. 67, 363 (1928).
- 3. BRIDGMAN P. W., Phys. Rev. 48, 893 (1935).
- 4. PIERMARINI G. J. and WEIR C. E., J. chem. Phys. 37, 1887 (1962).
- 5. WEIR C. E. and PIERMARINI G. J., J. Res. Nat. Bur. Standards 68A, 105 (1964).
- 6. JACOBS R. B., Phys. Rev. 54, 468 (1938).
- 7. JAMIESON J. C., J. Geol. 65, 334 (1957).
- ADAMS L. H. and DAVIS B. L., Proc. nat. Acad. Sci. 48, 982 (1962).
- 9. NAGASAKI H. and MINOMURA S., J. phys. Soc. Japan 19, 1496 (1964).
- KLEMENT W., Jr. and JAYARAMAN A., In Advances in Solid State Chemistry (Edited by H. Reiss) p. 289. Pergamon Press, Oxford, (1966).
- PISTORIUS C. W. F. T., J. Phys. Chem. Solids 25, 1477 (1964).
- 12. PISTORIUS C. W. F. T., ibid. 26, 1003 (1965).
- PISTORIUS C. W. F. T. and SYNMAN H. C., Z. phys. Chem. Neue Folge 43, 1 (1964).
- BRIDGMAN P. W., Physics of High Pressures. Bell, London (1958).
- KENNEDY G. C. and NEWTON R. C., In Solids Under Pressure (Edited by W. Paul and D. M. Warschauer), p. 163. McGraw-Hill, New York (1963).
- 16. JAYARAMAN A., KLEMENT W., NEWTON R. C. and KENNEDY G. C., J. Phys. Chem. Solids 24, 7 (1963).
- 17. JAYARAMAN A., KLEMENT W., Jr. and KENNEDY G. C., *Phys. Rev.* **130**, 2277 (1963).
- KENNEDY G. C. and LAMORI P. N., J. geophys. Res. 67, 851 (1962).
- KENNEDY G. C. and LAMORI P. N., In Progress in Very High Temperature Research (Edited by F. P. Bundy, W. R. Hibbard and H. M. Strong). Wiley, New York (1961).
- 20. LLOYD F. C., BECKETT C. W. and BOYD F. R., Jr., Science 164, 860 (1969).
- RAO C. N. R. and RAO K. J., In *Progress in Solid* State Chemistry (Edited by H. Reiss), Vol. 4, p. 131. Pergamon Press, Oxford (1966).
- 22. CLARK S. P., J. chem. Phys. 31, 1526 (1959).
- PISTORIUS C. W. F. T., J. Phys. Chem. Solids 26, 1543 (1965).

- 24. BRIDGMAN P. W., Proc. Am. Acad. Arts. Sci. 74, 21 (1945).
- 25. BRIDGMAN P. W., Proc. Am. Acad. Arts. Sci., 76, 1 (1945).
- 26. KNOF H. and MAISCH W. G., J. Phys. Chem. Solids 24, 1625 (1963).
- 27. MENARY J. W., UBBLOHDE A. R. and WOOD-WARD I., *Proc. R. Soc.* 208, 158 (1951).
- 28. BAUER S. H. and PORTER R. F., In Molten Salt

Chemistry (Edited by M. Blander), p. 623. Interscience, New York (1964).

- LARSON D. B. In *Physics of Solids at High Pressures* (Edited by C. T. Tomizuka and R. M. Emrick), p. 459. Academic Press, New York (1965).
- 30. GENSHAFT Yu. S., LIVSHITS L. D. and RYAB-ININ Yu. N., *Zh. Tek. Fiziki* 37, 179 (1967).
- 31. DANIELS W. B. and SKOULTCH1 A. 1., J. Phys. Chem. Solids 27, 1247 (1966).